

## The Green Function on a Compact 2D Manifold

MICHAEL TAYLOR

Let  $M$  be a compact 2D manifold, with a Riemannian metric tensor  $g$ . Say  $M$  has a coordinate system for which the components of this tensor are Hölder continuous, in  $C^r$ , for some  $r > 0$ . Let  $\Delta$  denote the Laplace-Beltrami operator, and  $\mathcal{G}$  the Green operator,

$$(1) \quad \begin{aligned} \mathcal{G} &= \Delta^{-1} && \text{on } V = \{u \in L^2(M) : \int u(x) dV(x) = 0\}, \\ &0 && \text{on } V^\perp. \end{aligned}$$

Let  $G(x, y)$  denote the integral kernel of  $\mathcal{G}$ , so

$$(2) \quad \mathcal{G}u(x) = \int_M G(x, y)u(y) dV(y).$$

Here and above,  $dV$  is the volume element on  $M$  associated to its metric tensor. One might think that if  $g$  were rough then  $G$  might be hard to analyze with any precision. Here we show that one can say a good bit about it, thanks to special structure in 2D. For simplicity we assume  $M$  is oriented, though we could remove that hypothesis.

First, note that  $g$  provides for  $M$  a canonical smooth structure. Indeed, under the hypotheses above,  $M$  has local isothermal coordinates (cf. [T], §III.9). These coordinates give  $M$  the structure of a Riemann surface, hence of a  $C^\infty$  manifold. In these coordinates, the components of  $g$  have the same regularity as hypothesized above. Furthermore, there exist *smooth* metric tensors on  $M$  in the same conformal class as  $g$  (including “canonical” metrics, with constant curvature). Pick one and denote it  $g_2$ .

Call the original metric tensor  $g_1$  and its Laplace operator, Green operator, and Green function  $\Delta_1, \mathcal{G}_1$ , and  $G_1(x, y)$ , and denote the parallel objects associated to  $g_2$  by  $\Delta_2, \mathcal{G}_2$ , and  $G_2(x, y)$ , so

$$(3) \quad \mathcal{G}_2u(x) = \int_M G_2(x, y)u(y) dV_2(y),$$

where  $dV_2$  is the volume element on  $M$  associated to  $g_2$ . Now  $G_2(x, y)$  is the integral kernel of a classical pseudodifferential operator of order  $-2$ ; it is an object with transparent structure. It remains to compare  $G_1(x, y)$  and  $G_2(x, y)$ .

Note that

$$(4) \quad \Delta_j = *_j d * d,$$

where  $*$  is the Hodge star on 1-forms and  $*_j$  takes 2-forms to 0-forms. Thus

$$(5) \quad \Delta_1 = A \Delta_2, \quad A = \frac{dV_2}{dV_1}.$$

If  $\int_M u(y) dV_1(y) = 0$ , then  $\int_M u(y) A(y)^{-1} dV_2(y) = 0$ , and

$$(6) \quad \mathcal{G}_1 u(x) = \mathcal{G}_2 A^{-1} u(x) \quad \text{mod const.}$$

Now

$$(7) \quad \begin{aligned} \mathcal{G}_2 A^{-1} u(x) &= \int_M G_2(x, y) A(y)^{-1} u(y) dV_2(y) \\ &= \int_M G_2(x, y) u(y) dV_1(y), \end{aligned}$$

so

$$(8) \quad \mathcal{G}_1 u(x) = \int_M G_2(x, y) u(y) dV_1(y) \quad \text{mod const.}, \quad \text{if } \int_M u(y) dV_1(y) = 0.$$

Of course

$$(9) \quad \mathcal{G}_1 u(x) = \int_M G_1(x, y) u(y) dV_1(y).$$

From here a brief calculation yields

$$(10) \quad G_1(x, y) = G_2(x, y) - \frac{\eta(x) + \eta(y)}{V_1(M)} + \frac{B}{V_1(M)},$$

where  $V_1(M) = \int_M 1 dV_1(M)$  and

$$(11) \quad \eta(x) = \int_M G_2(x, y) dV_1(y) = \mathcal{G}_2 \left( \frac{dV_1}{dV_2} \right) (x),$$

and

$$(12) \quad B = \frac{1}{V_1(M)} \int_M \eta(y) dV_1(y).$$

Recall we are trying to understand the behavior of  $G_1(x, y)$ , and as we have said the behavior of  $G_2(x, y)$  is well understood. Now if the components of  $g_1$  are of class  $C^r$  in local isothermal coordinates, we have

$$(13) \quad \frac{dV_1}{dV_2} \in C^r(M),$$

and hence, since the coefficients of  $\Delta_2$  are smooth,

$$(14) \quad \eta \in C^{r+2}(M).$$

Now let's say we are in the following situation. In some (possibly poorly chosen) coordinate system,  $g_1$  has components that are Hölder continuous and have one derivative in  $L^2$ . Then the Gauss curvature  $K$  is a well defined distribution. Assume

$$(15) \quad K \in L^4(M).$$

Now move to isothermal coordinates. Then the components of  $g_1$  belong to an  $L^4$ -Sobolev space:

$$(16) \quad H^{2,4} \subset C^{3/2}.$$

Hence, in isothermal coordinates, we have

$$(17) \quad \frac{dV_1}{dV_2} \in H^{2,4}$$

and hence (10) holds, with

$$(18) \quad \eta \in H^{4,4}(M) \subset C^{7/2}(M).$$

This is the degree of regularity one has in isothermal coordinates. In bad coordinates the regularity might be less.

EXAMPLE. Suppose  $M \subset \mathbb{R}^3$  is a surface smooth of class  $C^{1+r} \cap H^{2,p}$ ,  $0 < r < 1$ ,  $2 < p < \infty$ , i.e.,  $M$  is locally the graph of functions with that regularity. One can use such "graph coordinates," in which the metric tensor induced from the Euclidean metric on  $\mathbb{R}^3$  has components in  $C^r \cap H^{1,p}$ . Hence local isothermal coordinates exist. Also the Gauss map is smooth of class  $C^r \cap H^{1,p}$ , so

$$(19) \quad K \in L^q, \quad q = \frac{p}{2}.$$

Thus, in local isothermal coordinates the metric tensor has components in  $H^{2,q}$ , and we have (10) with  $\eta \in H^{4,q}$ , in isothermal coordinates. If  $p = 8$ , then (18) holds. Isothermal coordinates and graph coordinates are related by diffeomorphisms smooth of class  $C^{1+r} \cap H^{2,p}$ . Hence, if  $p = 8$ , one has  $\eta \in C^{7/2}(M)$  in isothermal coordinates, but only  $\eta \in C^{1+r}(M)$  in graph coordinates. Note that we have the same function  $\eta$  in both cases, just a finer differential structure on  $M$  in one case.

## Reference

[T] M. Taylor, Tools for PDE, Math. Surveys and Monogr. #81, AMS, Providence, RI, 2000.