## Integration of Differential Forms

The calculus of differential forms provides a convenient setting for integration on manifolds, as we will explain in this appendix, due to the efficient way it keeps track of changes of variables.

A $k$-form $\beta$ on an open set $\mathcal{O} \subset \mathbb{R}^{n}$ has the form

$$
\begin{equation*}
\beta=\sum_{j} b_{j}(x) d x_{j_{1}} \wedge \cdots \wedge d x_{j_{k}} . \tag{G.1}
\end{equation*}
$$

Here $j=\left(j_{1}, \ldots, j_{k}\right)$ is a $k$-multi-index. We write $\beta \in \Lambda^{k}(\mathcal{O})$. The wedge product used in (G.1) has the anti-commutative property

$$
\begin{equation*}
d x_{\ell} \wedge d x_{m}=-d x_{m} \wedge d x_{\ell} \tag{G.2}
\end{equation*}
$$

so that if $\sigma$ is a permutation of $\{1, \ldots, k\}$, we have

$$
\begin{equation*}
d x_{j_{1}} \wedge \cdots \wedge d x_{j_{k}}=(\operatorname{sgn} \sigma) d x_{j_{\sigma(1)}} \wedge \cdots \wedge d x_{j_{\sigma(k)}} \tag{G.3}
\end{equation*}
$$

In particular, an $n$-form $\alpha$ on $\Omega \subset \mathbb{R}^{n}$ can be written

$$
\begin{equation*}
\alpha=A(x) d x_{1} \wedge \cdots \wedge d x_{n} . \tag{G.4}
\end{equation*}
$$

If $A \in L^{1}(\mathcal{O}, d x)$, we write

$$
\begin{equation*}
\int_{\mathcal{O}} \alpha=\int_{\mathcal{O}} A(x) d x \tag{G.5}
\end{equation*}
$$

the right side being the usual Lebesgue integral, developed in Chapter 7.
Suppose now $\Omega \subset \mathbb{R}^{n}$ is open and there is a $C^{1}$ diffeomorphism $F: \Omega \rightarrow$ $\mathcal{O}$. We define the pull-back $F^{*} \beta$ of the $k$-form $\beta$ in (G.1) as

$$
\begin{equation*}
F^{*} \beta=\sum_{j} b_{j}(F(x))\left(F^{*} d x_{j_{1}}\right) \wedge \cdots \wedge\left(F^{*} d x_{j_{k}}\right) \tag{G.6}
\end{equation*}
$$

where

$$
\begin{equation*}
F^{*} d x_{j}=\sum_{\ell} \frac{\partial F_{j}}{\partial x_{\ell}} d x_{\ell} \tag{G.7}
\end{equation*}
$$

the algebraic computation in (G.6) being performed using the rule (G.3).
If $B=\left(b_{\ell m}\right)$ is an $n \times n$ matrix, then, by (G.3) and the formula for the determinant given in (7.77) (and (7.83)),

$$
\begin{align*}
\left(\sum_{m} b_{1 m} d x_{m}\right) & \wedge\left(\sum_{m} b_{2 m} d x_{m}\right) \wedge \cdots \wedge\left(\sum_{m} b_{n m} d x_{m}\right) \\
& =\left(\sum_{\sigma}(\operatorname{sgn} \sigma) b_{1 \sigma(1)} b_{2 \sigma(2)} \cdots b_{n \sigma(n)}\right) d x_{1} \wedge \cdots \wedge d x_{n}  \tag{G.8}\\
& =(\operatorname{det} B) d x_{1} \wedge \cdots \wedge d x_{n}
\end{align*}
$$

Hence, if $F: \Omega \rightarrow \mathcal{O}$ is a $C^{1}$ map and $\alpha$ is an $n$-form on $\mathcal{O}$, as in (G.4), then

$$
\begin{equation*}
F^{*} \alpha=\operatorname{det} D F(x) A(F(x)) d x_{1} \wedge \cdots \wedge d x_{n} \tag{G.9}
\end{equation*}
$$

This formula is especially significant in light of the change of variable formula

$$
\begin{equation*}
\int_{\mathcal{O}} A(x) d x=\int_{\Omega} A(F(x))|\operatorname{det} D F(x)| d x \tag{G.10}
\end{equation*}
$$

when $F: \Omega \rightarrow \mathcal{O}$ is a $C^{1}$ diffeomorphism, given in Theorem 7.2. The only difference between the right side of (G.10) and $\int_{\Omega} F^{*} \alpha$ is the absolute value sign around $\operatorname{det} D F(x)$. We say a $C^{1}$ map $F: \Omega \rightarrow \mathcal{O}$ is orientation preserving when $\operatorname{det} D F(x)>0$ for all $x \in \Omega$. In such a case, Theorem 7.2 yields

Proposition G.1. If $F: \Omega \rightarrow \mathcal{O}$ is a $C^{1}$ orientation-preserving diffeomorphism and $\alpha$ an integrable $n$-form on $\mathcal{O}$, then

$$
\begin{equation*}
\int_{\mathcal{O}} \alpha=\int_{\Omega} F^{*} \alpha . \tag{G.11}
\end{equation*}
$$

In Appendix H we will present another proof of the change of variable formula, making direct use of basic results on differential forms developed in this appendix.

In addition to the pull-back, there are some other operations on differential forms. The wedge product of $d x_{\ell}$ 's extends to a wedge product on forms as follows. If $\beta \in \Lambda^{k}(\mathcal{O})$ has the form (G.1) and if

$$
\begin{equation*}
\alpha=\sum_{i} a_{i}(x) d x_{i_{1}} \wedge \cdots \wedge d x_{i_{\ell}} \in \Lambda^{\ell}(\mathcal{O}), \tag{G.12}
\end{equation*}
$$

define

$$
\begin{equation*}
\alpha \wedge \beta=\sum_{i, j} a_{i}(x) b_{j}(x) d x_{i_{1}} \wedge \cdots \wedge d x_{i_{\ell}} \wedge d x_{j_{1}} \wedge \cdots \wedge d x_{j_{k}} \tag{G.13}
\end{equation*}
$$

in $\Lambda^{k+\ell}(\mathcal{O})$. We retain the equivalences (G.3). It follows that

$$
\begin{equation*}
\alpha \wedge \beta=(-1)^{k \ell} \beta \wedge \alpha \tag{G.14}
\end{equation*}
$$

It is also readily verified that

$$
\begin{equation*}
F^{*}(\alpha \wedge \beta)=\left(F^{*} \alpha\right) \wedge\left(F^{*} \beta\right) . \tag{G.15}
\end{equation*}
$$

Another important operator on forms is the exterior derivative:

$$
\begin{equation*}
d: \Lambda^{k}(\mathcal{O}) \longrightarrow \Lambda^{k+1}(\mathcal{O}) \tag{G.16}
\end{equation*}
$$

defined as follows. If $\beta \in \Lambda^{k}(\mathcal{O})$ is given by (G.1), then

$$
\begin{equation*}
d \beta=\sum_{j, \ell} \frac{\partial b_{j}}{\partial x_{\ell}} d x_{\ell} \wedge d x_{j_{1}} \wedge \cdots \wedge d x_{j_{k}} \tag{G.17}
\end{equation*}
$$

The antisymmetry $d x_{m} \wedge d x_{\ell}=-d x_{\ell} \wedge d x_{m}$, together with the identity $\partial^{2} b_{j} / \partial x_{\ell} \partial x_{m}=\partial^{2} b_{j} / \partial x_{m} \partial x_{\ell}$, implies

$$
\begin{equation*}
d(d \beta)=0, \tag{G.18}
\end{equation*}
$$

for any smooth differential form $\beta$. We also have a product rule:

$$
\begin{equation*}
d(\alpha \wedge \beta)=(d \alpha) \wedge \beta+(-1)^{j} \alpha \wedge(d \beta), \quad \alpha \in \Lambda^{j}(\mathcal{O}), \beta \in \Lambda^{k}(\mathcal{O}) \tag{G.19}
\end{equation*}
$$

The exterior derivative has the following important property under pullbacks:

$$
\begin{equation*}
F^{*}(d \beta)=d F^{*} \beta, \tag{G.20}
\end{equation*}
$$

if $\beta \in \Lambda^{k}(\mathcal{O})$ and $F: \Omega \rightarrow \mathcal{O}$ is a smooth map. To see this, extending (G.19) to a formula for $d\left(\alpha \wedge \beta_{1} \wedge \cdots \wedge \beta_{\ell}\right)$ and using this to apply $d$ to $F^{*} \beta$, we have

$$
\begin{align*}
d F^{*} \beta= & \sum_{j, \ell} \frac{\partial}{\partial x_{\ell}}\left(b_{j} \circ F(x)\right) d x_{\ell} \wedge\left(F^{*} d x_{j_{1}}\right) \wedge \cdots \wedge\left(F^{*} d x_{j_{k}}\right)  \tag{G.21}\\
& +\sum_{j, \nu}( \pm) b_{j}(F(x))\left(F^{*} d x_{j_{1}}\right) \wedge \cdots \wedge d\left(F^{*} d x_{j_{\nu}}\right) \wedge \cdots \wedge\left(F^{*} d x_{j_{k}}\right)
\end{align*}
$$

Now the definition (G.6)-(G.7) of pull-back gives directly that

$$
\begin{equation*}
F^{*} d x_{i}=\sum_{\ell} \frac{\partial F_{i}}{\partial x_{\ell}} d x_{\ell}=d F_{i}, \tag{G.22}
\end{equation*}
$$

and hence $d\left(F^{*} d x_{i}\right)=d d F_{i}=0$, so only the first sum in (G.21) contributes to $d F^{*} \beta$. Meanwhile,

$$
\begin{equation*}
F^{*} d \beta=\sum_{j, m} \frac{\partial b_{j}}{\partial x_{m}}(F(x))\left(F^{*} d x_{m}\right) \wedge\left(F^{*} d x_{j_{1}}\right) \wedge \cdots \wedge\left(F^{*} d x_{j_{k}}\right), \tag{G.23}
\end{equation*}
$$

so (G.20) follows from the identity

$$
\sum_{\ell} \frac{\partial}{\partial x_{\ell}}\left(b_{j} \circ F(x)\right) d x_{\ell}=\sum_{m} \frac{\partial b_{j}}{\partial x_{m}}(F(x)) F^{*} d x_{m}
$$

which in turn follows from the chain rule.
Here is another important consequence of the chain rule. Suppose $F$ : $\Omega \rightarrow \mathcal{O}$ and $\psi: \mathcal{O} \rightarrow U$ are smooth maps between open subsets of $\mathbb{R}^{n}$. We claim that for any form $\alpha$ of any degree,

$$
\begin{equation*}
\psi \circ F=\varphi \Longrightarrow \varphi^{*} \alpha=F^{*} \psi^{*} \alpha . \tag{G.24}
\end{equation*}
$$

It suffices to check (G.24) for $\alpha=d x_{j}$. Then (G.7) gives the basic identity $\psi^{*} d x_{j}=\sum\left(\partial \psi_{j} / \partial x_{\ell}\right) d x_{\ell}$. Consequently,

$$
\begin{equation*}
F^{*} \psi^{*} d x_{j}=\sum_{\ell, m} \frac{\partial F_{\ell}}{\partial x_{m}} \frac{\partial \psi_{j}}{\partial x_{\ell}} d x_{m}, \quad \varphi^{*} d x_{j}=\sum_{m} \frac{\partial \varphi_{j}}{\partial x_{m}} d x_{m} \tag{G.25}
\end{equation*}
$$

but the identity of these forms follows from the chain rule:

$$
\begin{equation*}
D \varphi=(D \psi)(D F) \Longrightarrow \frac{\partial \varphi_{j}}{\partial x_{m}}=\sum_{\ell} \frac{\partial \psi_{j}}{\partial x_{\ell}} \frac{\partial F_{\ell}}{\partial x_{m}} \tag{G.26}
\end{equation*}
$$

One can define a $k$-form on an $n$-dimensional manifold $M$ as follows. Say $M$ is covered by open sets $\mathcal{O}_{j}$ and there are coordinate charts $F_{j}: \Omega_{j} \rightarrow \mathcal{O}_{j}$, with $\Omega_{j} \subset \mathbb{R}^{n}$ open. A collection of forms $\beta_{j} \in \Lambda^{k}\left(\Omega_{j}\right)$ is said to define a $k$-form on $M$ provided the following compatibility condition holds. If $\mathcal{O}_{i} \cap \mathcal{O}_{j} \neq \emptyset$ and we consider $\Omega_{i j}=F_{i}^{-1}\left(\mathcal{O}_{i} \cap \mathcal{O}_{j}\right)$ and diffeomorphisms

$$
\begin{equation*}
\varphi_{i j}=F_{j}^{-1} \circ F_{i}: \Omega_{i j} \longrightarrow \Omega_{j i} \tag{G.27}
\end{equation*}
$$

we require

$$
\begin{equation*}
\varphi_{i j}^{*} \beta_{j}=\beta_{i} \tag{G.28}
\end{equation*}
$$

The fact that this is a consistent definition is a consequence of (G.24). For example, if $G: M \rightarrow R^{m}$ is a smooth map and $\gamma$ is a $k$-form on $\mathbb{R}^{m}$, then there is a well-defined $k$-form $\beta=G^{*} \gamma$ on $M$, represented in such coordinate charts by $\beta_{j}=\left(G \circ F_{j}\right)^{*} \gamma$. Similarly, if $\beta$ is a $k$-form on $M$ as defined above and $G: U \rightarrow M$ is smooth, with $U \subset \mathbb{R}^{m}$ open, then $G^{*} \beta$ is a well-defined $k$-form on $U$.

We give an intrinsic definition of $\int_{M} \alpha$ when $\alpha$ is an $n$-form on $M$, provided $M$ is oriented, i.e., there is a coordinate cover as above such that $\operatorname{det} D \varphi_{j k}>0$. The object called an "orientation" on $M$ can be identified as an equivalence class of nowhere vanishing $n$-forms on $M$, two such forms being equivalent if one is a multiple of another by a positive function in $C^{\infty}(\Omega)$. A member of this equivalence class, say $\omega$, defines the orientation. The standard orientation on $\mathbb{R}^{n}$ is determined by $d x_{1} \wedge \cdots \wedge d x_{n}$. The equivalence class of positive multiples $a(x) \omega$ is said to consist of "positive" forms. A smooth map $\psi: S \rightarrow M$ between oriented $n$-dimensional manifolds preserves orientation provided $\psi^{*} \sigma$ is positive on $S$ whenever $\sigma \in \Lambda^{n}(M)$ is positive. We mention that there exist surfaces that cannot be oriented, such as the famous "Möbius strip."

We define the integral of an $n$-form over an oriented $n$-dimensional manifold as follows. First, if $\alpha$ is an $n$-form supported on an open set $\mathcal{O} \subset \mathbb{R}^{n}$, given by (G.4), then we define $\int_{\mathcal{O}} \alpha$ by (G.5).

More generally, if $M$ is an $n$-dimensional manifold with an orientation, say the image of an open set $\mathcal{O} \subset \mathbb{R}^{n}$ by $\varphi: \mathcal{O} \rightarrow M$, carrying the natural orientation of $\mathcal{O}$, we can set

$$
\begin{equation*}
\int_{M} \alpha=\int_{\mathcal{O}} \varphi^{*} \alpha \tag{G.29}
\end{equation*}
$$

for an $n$-form $\alpha$ on $M$. If it takes several coordinate patches to cover $M$, define $\int_{M} \alpha$ by writing $\alpha$ as a sum of forms, each supported on one patch.

We need to show that this definition of $\int_{M} \alpha$ is independent of the choice of coordinate system on $M$ (as long as the orientation of $M$ is respected). Thus, suppose $\varphi: \mathcal{O} \rightarrow U \subset M$ and $\psi: \Omega \rightarrow U \subset M$ are both coordinate patches, so that $F=\psi^{-1} \circ \varphi: \mathcal{O} \rightarrow \Omega$ is an orientation-preserving diffeomorphism. We need to check that, if $\alpha$ is an $n$-form on $M$, supported on $U$, then

$$
\begin{equation*}
\int_{\mathcal{O}} \varphi^{*} \alpha=\int_{\Omega} \psi^{*} \alpha \tag{G.30}
\end{equation*}
$$

To establish this, we use (G.24). This implies that the left side of (G.30) is equal to

$$
\begin{equation*}
\int_{\mathcal{O}} F^{*}\left(\psi^{*} \alpha\right) \tag{G.31}
\end{equation*}
$$

which is equal to the right side of (G.30), by (G.11) (with slightly altered notation). Thus the integral of an $n$-form over an oriented $n$-dimensional manifold is well defined.

We turn now to the Gauss-Green-Stokes formula for differential forms, commonly called simply the Stokes formula. This involves integrating a $k$-form over a $k$-dimensional manifold with boundary. We first define that concept. Let $S$ be a smooth $k$-dimensional manifold, and let $M$ be an open subset of $S$, such that its closure $\bar{M}\left(\right.$ in $\left.\mathbb{R}^{N}\right)$ is contained in $S$. Its boundary is $\partial M=\bar{M} \backslash M$. We say $\bar{M}$ is a smooth surface with boundary if also $\partial M$ is a smooth $(k-1)$-dimensional surface. In such a case, any $p \in \partial M$ has a neighborhood $U \subset S$ with a coordinate chart $\varphi: \mathcal{O} \rightarrow U$, where $\mathcal{O}$ is an open neighborhood of 0 in $\mathbb{R}^{k}$, such that $\varphi(0)=p$ and $\varphi \operatorname{maps}\left\{x \in \mathcal{O}: x_{1}=0\right\}$ onto $U \cap \partial M$.

If $S$ is oriented, then $\bar{M}$ is oriented, and $\partial M$ inherits an orientation, uniquely determined by the following requirement: if

$$
\begin{equation*}
\bar{M}=\mathbb{R}_{-}^{k}=\left\{x \in \mathbb{R}^{k}: x_{1} \leq 0\right\} \tag{G.32}
\end{equation*}
$$

then $\partial M=\left\{\left(x_{2}, \ldots, x_{k}\right)\right\}$ has the orientation determined by $d x_{2} \wedge \cdots \wedge d x_{k}$.
We can now state the Stokes formula.
Proposition G.2. Given a compactly supported $(k-1)$-form $\beta$ of class $C^{1}$ on an oriented $k$-dimensional surface $\bar{M}$ (of class $C^{2}$ ) with boundary $\partial M$, with its natural orientation,

$$
\begin{equation*}
\int_{M} d \beta=\int_{\partial M} \beta \tag{G.33}
\end{equation*}
$$

Proof. Using a partition of unity and invariance of the integral and the exterior derivative under coordinate transformations, it suffices to prove this when $\bar{M}$ has the form (G.32). In that case, we will be able to deduce (G.33) from the Fundamental Theorem of Calculus. Indeed, if

$$
\begin{equation*}
\beta=b_{j}(x) d x_{1} \wedge \cdots \wedge \widehat{d x}_{j} \wedge \cdots \wedge d x_{k} \tag{G.34}
\end{equation*}
$$

with $b_{j}(x)$ of bounded support, we have

$$
\begin{equation*}
d \beta=(-1)^{j-1} \frac{\partial b_{j}}{\partial x_{j}} d x_{1} \wedge \cdots \wedge d x_{k} \tag{G.35}
\end{equation*}
$$

If $j>1$, we have

$$
\begin{equation*}
\int_{M} d \beta=(-1)^{j-1} \int\left\{\int_{-\infty}^{\infty} \frac{\partial b_{j}}{\partial x_{j}} d x_{j}\right\} d x^{\prime}=0 \tag{G.36}
\end{equation*}
$$

and also $\kappa^{*} \beta=0$, where $\kappa: \partial M \rightarrow \bar{M}$ is the inclusion. On the other hand, for $j=1$, we have

$$
\begin{align*}
\int_{M} d \beta & =\int\left\{\int_{-\infty}^{0} \frac{\partial b_{1}}{\partial x_{1}} d x_{1}\right\} d x_{2} \cdots d x_{k} \\
& =\int b_{1}\left(0, x^{\prime}\right) d x^{\prime}  \tag{G.37}\\
& =\int_{\partial M} \beta
\end{align*}
$$

This proves Stokes' formula (G.33).
The reason we required $\bar{M}$ to be a surface of class $C^{2}$ in Proposition G. 2 is the following. Due to the formulas (G.6)-(G.7) for a pull-back, if $\beta$ is of class $C^{j}$ and $F$ is of class $C^{\ell}$, then $F^{*} \beta$ is generally of class $C^{\mu}$, with $\mu=\min (j, \ell-1)$. Thus, if $j=\ell=1, F^{*} \beta$ might be only of class $C^{0}$, so there is not a well-defined notion of a differential form of class $C^{1}$ on a $C^{1}$ surface, though such a notion is well defined on a $C^{2}$ surface. This problem can be overcome, and one can extend Proposition G. 2 to the case where $\bar{M}$ is a $C^{1}$ surface and $\beta$ is a $(k-1)$-form with the property that both $\beta$ and $d \beta$ are continuous. One can go further and formulate (G.33) for a ( $k-1$ )-form $\beta$ with the property that

$$
\begin{equation*}
\beta, d \beta \in L^{\infty}(M), \quad \iota^{*} \beta \in L^{\infty}(\partial M), \tag{G.38}
\end{equation*}
$$

where $\iota: \partial M \rightarrow \bar{M}$ is the natural inclusion, a class of forms that can be shown to be invariant under bi-Lipschitz maps. (It can be shown that the first two conditions in (G.38) imply $\iota^{*} \beta \in H^{1,1}(\partial M)^{\prime}$.) We will not go into the details. However, in Appendix I we will present an elementary treatment of (G.33), stated in a more classical language, when $M$ is an open domain in $\mathbb{R}^{k}$ whose boundary is locally the graph of a Lipschitz function. A far reaching extension, due to H. Federer, can be found in [Fed]; see also [EG].

The calculus of differential forms has many applications to differential equations, differential geometry, and topology. More on this can be found in $[\mathbf{S p i}]$ and also in $[\mathbf{T} 1]$ (particularly Chapters 1,5 , and 10). To end this appendix, we make use of the calculus of differential forms to provide simple proofs of some important topological results of Brouwer. The first two results concern retractions. If $Y$ is a subset of $X$, by definition a retraction of $X$ onto $Y$ is a map $\varphi: X \rightarrow Y$ such that $\varphi(x)=x$ for all $x \in Y$.

Proposition G.3. There is no smooth retraction $\varphi: B \rightarrow S^{n-1}$ of the closed unit ball $B$ in $\mathbb{R}^{n}$ onto its boundary $S^{n-1}$.

In fact, it is just as easy to prove the following more general result. The approach we use is adapted from [Kan].

Proposition G.4. If $\bar{M}$ is a compact oriented $n$-dimensional manifold with nonempty boundary $\partial M$, there is no smooth retraction $\varphi: \bar{M} \rightarrow \partial M$.

Proof. You can pick $\omega \in \Lambda^{n-1}(\partial M)$ to be an $(n-1)$-form on $\partial M$ such that $\int_{\partial M} \omega>0$. Now apply Stokes' Theorem to $\beta=\varphi^{*} \omega$. If $\varphi$ is a retraction, then $\varphi \circ j(x)=x$, where $j: \partial M \hookrightarrow \bar{M}$ is the natural inclusion. Hence $j^{*} \varphi^{*} \omega=\omega$, so we have

$$
\begin{equation*}
\int_{\partial M} \omega=\int_{M} d \varphi^{*} \omega . \tag{G.39}
\end{equation*}
$$

But $d \varphi^{*} \omega=\varphi^{*} d \omega=0$, so the integral (G.39) is zero. This is a contradiction, so there can be no retraction.

A simple consequence of this is the famous Brouwer Fixed-Point Theorem.

Theorem G.5. If $F: B \rightarrow B$ is a continuous map on the closed unit ball in $\mathbb{R}^{n}$, then $F$ has a fixed point.

Proof. First, an approximation argument shows that if there is a continuous such $F$ without a fixed point, then there is a smooth one, so assume $F$ : $B \rightarrow B$ is smooth. We are claiming that $F(x)=x$ for some $x \in B$. If not,
then for each $x \in B$ define $\varphi(x)$ to be the endpoint of the ray from $F(x)$ to $x$, continued until it hits $\partial B=S^{n-1}$. It is clear that $\varphi$ would be a smooth retraction, contradicting Proposition G.3.

Remark. Typical proofs of the Brouwer Fixed-Point Theorem use concepts of algebraic topology; cf. [Spa]. In fact, the proof of Proposition G. 4 contains a germ of de Rham cohomology. See [T1], Chapter $1, \S 19$ for more on this.

An integral calculus proof of the Brouwer Fixed-Point Theorem that does not involve differential forms is given in [DS], Vol. 1, pp. 467-470. One might compare it with the proof given above.

